



## Metal-Decorated Silicene for Hydrogen Storage

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**Abstract:** Environmental and energy security concerns makes renewable energy research imperative. A key issue that needs to be addressed before the implementation of a renewable energy system is the issue of energy storage, in which materials-based hydrogen storage is considered as an option. In this paper, the metal-decorated silicene was studied for hydrogen storage applications. These were performed using the Vienna Ab Initio Simulation Package (VASP), a plane wave-based code. The potential was described using the projector-augmented wave (PAW) method. The exchange-correlation energy was approximated using Generalized Gradient Approximation (GGA) functional in the Perdew, Burke, and Ernzerhof (PBE) form. The k points were gamma-centered and automatically generated using the Monkhorst-Pack scheme. The energy cutoff, number of k points, vacuum spacing, silicene lattice constant, and silicene buckling height were determined using various convergence tests. The transition metals Co, Ni, Pd, and Pt were used for the metal decoration. Relaxation calculations were done to determine the most stable adsorption site for the metal adatoms. Results show the hollow site to be the most stable among the four adsorption sites (bridge, hollow, top, and valley). Static calculations with Pd and Pt decoration also confirmed this.

**Key Words:** Hydrogen Storage; Silicene; Metal Decoration, Adsorption

### 1. INTRODUCTION

The continuous increase in world energy demand and the negative impact of fossil fuel sources in the environment creates a need for diversifying energy sources (Höök & Tang, 2013; IEA, 2014). The most preferable alternative is a renewable energy source which produces little or no emissions such as solar or wind power. However, the power output of these sources may become unstable because of unfavorable weather conditions such as cloudy weather or low wind speed. Hence, there is a need for an effective energy storage method so that renewable

systems can provide a steady supply of energy even during these periods (Converse, 2006). One of the alternatives for effective energy storage is through the use of hydrogen, an energy carrier.

As an energy carrier, hydrogen can be used as energy source by conversion to heat or electricity via combustion or electrochemical reactions (Pudukudy, Yaakob, Mohammad, Narayanan, & Sopian, 2014). Its combustion produces water and no harmful emissions except NO<sub>x</sub> compounds (Cipriani et al., 2014; Midilli & Dincer, 2008; Momirlan & Veziroglu, 2005; Najjar, 2013; Pudukudy et al., 2014; Sharaf & Orhan, 2014) which can be remedied by modifying the engine design (Momirlan & Veziroglu,

2005). This makes hydrogen a suitable energy carrier for a greener energy system

Hydrogen has 143 MJ·kg<sup>-1</sup> energy weight density—about three times the weight density of liquid hydrocarbon fuels. However, gaseous hydrogen has low energy volumetric density of 0.0107 MJ·L<sup>-1</sup>. Physical storage through liquid hydrogen (LH<sub>2</sub>) is also expensive (Najjar, 2013) and has problems with leakage (Cipriani et al., 2014; Najjar, 2013), which raises issues on the safety of LH<sub>2</sub> storage.

An alternative to physical storage is materials-based hydrogen storage through chemical hydrides and adsorption materials (Cipriani et al., 2014; Durbin & Malardier-Jugroot, 2013; Najjar, 2013; Pudukudy et al., 2014), which is relatively safer (Pudukudy et al., 2014). One of the popular materials actively studied in this field are carbon nanomaterials (Pudukudy et al., 2014), where many hydrogen storage studies has been conducted (Gopalsamy & Subramanian, 2014; Ströbel, Garche, Moseley, Jörissen, & Wolf, 2006; Wang et al., 2013a; Yang, Jung, Kim, & Park, 2012; Züttel, Wenger, Sudan, Mauron, & Orimo, 2004). Examples of nanomaterials include nanotubes, nanoflowers, nanobamboos, graphene, and many others. Calculations for hydrogen storage capacities give of 3 to 14 wt% storage, but experiments currently yield 0 to 3 wt % storage (Ströbel et al., 2006).

Many studies on CNTs and graphene showed that metal adsorption and doping could improve the hydrogen storage capacity of carbon nanomaterials (Ströbel et al., 2006). Metals that were used for improving the hydrogen storage capacity of graphene include alkali metals, alkaline earth metals, and transition metals (Chen et al., 2013; Choi, Min, Gwak, Paek, & Oh, 2014; Chu, Hu, Hu, Yang, & Deng, 2011; Gao et al., 2012; Guo et al., 2013; Kim, Lee, Hwang, Yun, & Chung, 2014; Li et al., 2012; Ma, Zhang, Xu, & Ji, 2014; Nachimuthu, Lai, & Jiang, 2014; Sigal, Rojas, & Leiva, 2011; Yadav, Zhu, & Singh, 2014; Zhou et al., 2014). Hydrogen molecules were found to have Coulombic or Kubas interaction with these metal dopants, whose binding energies allow for relatively high hydrogen storage (Wang et al., 2013a). In another mechanism called “spillover effect”, transition metals act as catalyst for hydrogen dissociation. The dissociated hydrogen atoms then forms bonds with carbon atoms and is said to have “spilled over” the carbon nanomaterial (Ströbel et al., 2006; Wang et al., 2013b). However, one problem is that metal adatoms tend to cluster on graphene, which significantly lowers the storage capacity. Calculations showed that this can be remedied by substitutional doping of graphene with nitrogen or boron (Nachimuthu et al., 2014).

Recently, the silicon analog of graphene,

silicene, was fabricated (Vogt et al., 2012). Metal adsorption studies showed stronger binding of metal adatoms with silicene as compared with graphene (Lin & Ni, 2012). Moreover, the metal-silicene binding energies are stronger than the cohesive energies of these metals. This means that silicene will not encounter the same metal clustering problems as graphene. Therefore, it is possible that silicene could also be used for hydrogen storage like graphene, since there are similar properties between the two. Also, it is interesting to study how silicene would be different from graphene in terms of hydrogen interaction, since silicene, unlike graphene, has higher lattice constant along the planar lattice vectors and is not perfectly flat but has a buckled structure (Lin & Ni, 2012).

## 2. METHODOLOGY

Spin-polarized ground state density functional theory (DFT) calculations were done using the Vienna Ab-Initio Simulation Package (VASP), a plane wave-based code (Kresse & Furthmüller, 1996a, 1996b; Kresse & Hafner, 1993, 1994). The projector-augmented wave (PAW) method was used to describe the potentials (Blöchl, 1994; Kresse & Joubert, 1999). The Generalized Gradient Approximation (GGA) functional in the Perdew, Burke, and Ernzerhof (PBE) form was used to approximate the exchange-correlation energy (Perdew, Burke, & Ernzerhof, 1996, 1997). The k point mesh was automatically generated using the Monkhorst Pack scheme. Gaussian smearing with broadening width of 0.1 eV was used. For static calculations, the energy convergence criterion of 1x10<sup>-4</sup> eV was used. The atoms were relaxed until the atomic forces were below 10<sup>-2</sup> eV/Å. For convergence test calculations involving one to two atoms, only the gamma point and default plane wave energy cutoff of metal was used. The static calculations to determine the silicene buckling height and lattice constant used an 11x11x11 gamma-centered k point mesh with a plane wave energy cutoff of 500 eV. Relaxation and static calculations used 7x7x1 k point mesh and 400 eV energy cutoff. The transition metals Co, Ni, Pd, and Pt were used for the metal decoration. The binding energy of the metal adatom with the silicene substrate was computed using the equation:

$$E_B = E_{S-M} - (E_S + E_M) \quad (\text{Eq. 1})$$

where:

$$E_B = \text{Binding energy of metal adatom with silicene}$$

$$E_{S-M} = \text{Energy of metal-decorated silicene}$$

$E_S$  = Energy of pristine silicene  
 $E_M$  = Energy of isolated metal adatom

Four adsorption sites were tested: the bridge site (B), hollow site (H), top site of the higher sublattice plane (T), and the valley site or the top site of the lower sublattice plane (V). Figure 1 shows the location of these four sites on a 4x4 silicene cell.

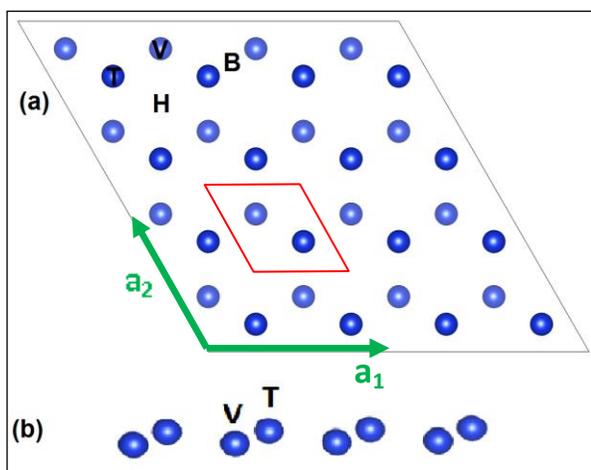


Fig. 1. (a) A 4x4 silicene cell showing the four possible adsorption sites. The primitive unit cell is enclosed in red box. This primitive cell is translated along the  $a_1$  and  $a_2$  lattice vectors to form the 4x4 silicene cell. (b) A side view of 4x4 silicene shows the distinction between the top site (T) and valley site (V).

For convergence tests involving isolated atoms, the cubic cell was used. For convergence test of silicene buckling height and lattice constant using static calculations, a 1x1 primitive rhombohedral cell was used (Fig. 1). The buckling height and lattice constant from these calculations were implemented in the relaxation calculations (which utilized a 4x4 unit cell) for finding the most stable metal adsorption site on silicene. A vacuum height of 18 Å was used for these calculations and dipole correction along the z-axis was implemented.

The results from VASP were visualized using the VESTA (Visualization for Electronic and Structural Analysis) software (Momma & Izumi, 2008, 2011). VESTA was also used to measure the distances between metal adsorbate and silicon atoms in the relaxed structure.

### 3. RESULTS AND DISCUSSION

#### 3.1 Convergence Tests

Convergence tests were done to verify the optimal k point sampling mesh and plane wave energy cutoff of the system. For the k point sampling convergence test, the 4 x 4 pristine silicene was used since this system has larger dimensions than the isolated metal adatom. For the plane-wave energy cutoff convergence test, the metal adatoms were used since these elements required larger energy cutoff.

##### 3.1.1 K point convergence test

The number of k points along the  $a_1$  and  $a_2$  lattice vectors ( $n_k$ ) was increased from 1 to 12 with increment of one. The absolute value of the energy differences ( $\Delta E$ ) per increase in  $n_k$  was computed. Figure 2 shows the results. The consistent decrease in  $\Delta E$  can be seen from the Fig. 2 (inset) starting from  $n_k=7$ . However, the energy difference that is comparable with 1meV was attained at  $n_k=8$ .

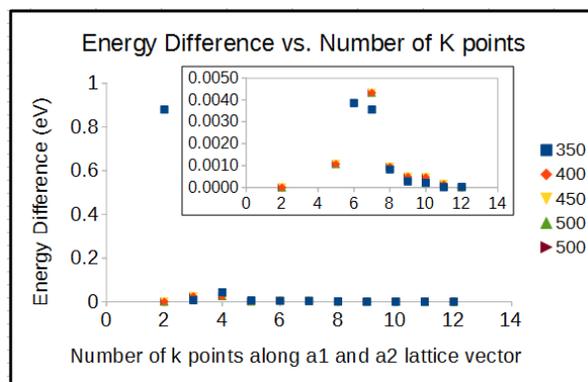


Fig. 2. The resulting energy differences vs. number of k points

##### 3.1.2 Plane wave energy cutoff convergence test

The plane wave energy cutoff ( $E_{cut}$ ) was varied from 150 to 500 eV with increment of 50 eV. The energy difference of  $\Delta E$  was also computed. Figure 3 shows the results of this convergence test. A steady decrease in  $\Delta E$  was observed as the  $E_{cut}$  was increased. The difference of 1meV was observed for  $E_{cut}=450$ eV, except for Pd, which required  $E_{cut}=500$ eV for an  $\Delta E =1$ meV to be attained (Fig. 3, inset). The computed  $\Delta E \sim 2$ meV for Pd at  $E_{cut}=450$ eV can still be

considered to be an acceptable energy difference to attain accurate calculations.

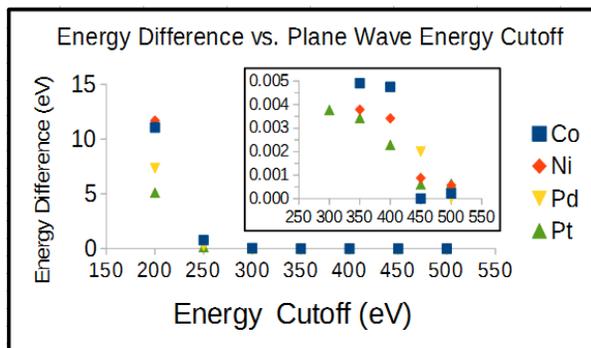


Fig. 3. The resulting energy differences vs. plane wave energy cutoff.

### 3.2 Determining the buckling height and lattice constant of silicene

The buckling height ( $\Delta z$ ) and lattice constant ( $a$ ) along the  $a_1$  and  $a_2$  directions were determined using static calculations. The lattice constant was varied from 3.78 Å to 4.00 Å, with increment of 0.01 Å, while  $\Delta z$  was varied from 0.40 Å to 0.54 Å, with increment of 0.01 Å. The energy from different combinations of these parameters was calculated. The energy of a flat silicene ( $\Delta z=0.00$  Å) with varying lattice constant was also calculated. The parameter combination that yielded the lowest energy is considered to be the most stable and the one that was used for the succeeding calculations. Figure 4 shows how the energy varies for different lattice constant and buckling height values. It can be easily seen that the flat silicene is the most unstable among the different systems (Fig. 4, inset). The figure shows that very close minima are obtained from different  $\Delta z$  –  $a$  combinations. A closer examination shows that the lowest energy was obtained for  $a=3.87$  Å and  $\Delta z=0.45$  Å. Hence, these values were used for the relaxation calculations (Sec. 3.3) in determining the optimal binding site for the metal adsorbate.

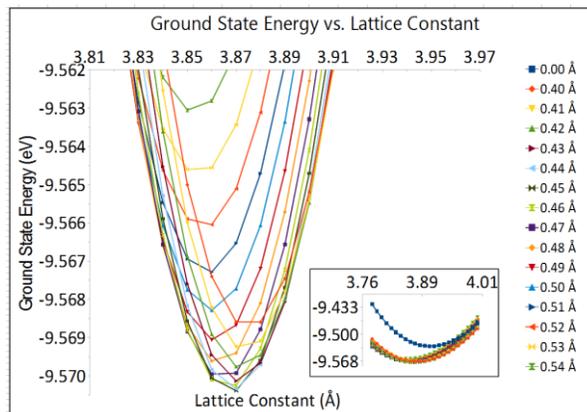


Fig. 4. The ground state energy of silicene for different lattice parameter and buckling height combinations.

### 3.3 Verifying the most stable adatom site

Both relaxation and static calculations were used to verify the most stable metal adsorption site. For the relaxation calculations, the metal and silicon atoms were allowed to move but the cell size and shape were maintained constant. The adsorption site with the highest binding energy ( $E_b$ ) is considered to be the most stable based from our definition in Eq. 1. For the static calculations, the silicon atoms are fixed and the height of metal adatom above the binding site ( $z$ ) was varied. The calculated ground state energies for different heights and adsorption sites were compared. The minimum point is considered to be the most stable site for the metal adsorbate.

#### 3.3.1 Relaxation Calculations

The computed  $E_b$  was used to calculate the ratio of the binding energy of metal adsorbate with silicene over the cohesive energy of the metal bulk ( $E_b/E_c$ ). The standard value of  $E_c$  was used from Kittel (1996). A higher  $E_b/E_c$  ratio means that there will be lower possibility of adatom clustering, which would result in better  $H_2$  uptake by the material. Table 1 shows the computed  $E_b/E_c$  ratio. The  $E_b/E_c$  results from the study of Lin & Ni (2012) was compared with our computed value. The distance of metal adatom with the nearest silicon atom ( $d_{M-Si}$ ) was also obtained. Table 2 shows the computed  $d_{M-Si}$  values and the comparison with Lin and Ni's (2012)  $d_{M-Si}$  values.

Table 1. Binding Energy over Cohesive Energy Ratio for Metal-Decorated Silicene

Metal Adatom	Adsorption Site	$E_b/E_c$ Calculated	$E_b/E_c$ (Lin & Ni, 2012)	Diff
Co	B	0.865	-	-
	H	0.968	1.017	0.048
	T	0.835	0.896	0.061
	V	0.865	0.924	0.059
Ni	B	0.934	-	-
	H	1.070	1.076	0.006
	T	0.889	0.893	0.004
Pd	V	0.934	0.939	0.005
	B	0.965	-	-
	H	1.095	1.080	0.015
Pt	T	0.910	0.898	0.012
	V	0.965	0.951	0.014
	B	0.897	-	-
Pt	H	1.004	1.005	0.001
	T	0.875	0.878	0.003
	V	0.897	0.899	0.002

Table 2. Smallest Metal-Silicon Bond Length ( $d_{M-Si}$ )

Metal Adatom	Adsorption Site	$d_{M-Si}$ Calculated	$d_{M-Si}$ (Lin & Ni, 2012)	Diff
Co	B	2.27	-	-
	H	2.30	2.27	0.03
	T	2.26	2.29	0.03
	V	2.27	2.29	0.02
Ni	B	2.25	-	-
	H	2.32	2.31	0.01
	T	2.24	2.24	0.00
Pd	V	2.25	2.25	0.00
	B	2.42	-	-
	H	2.46	2.46	0.00
Pt	T	2.41	2.42	0.01
	V	2.42	2.42	0.00
	B	2.41	-	-
	H	2.43	2.42	0.01
Pt	T	2.42	2.43	0.01
	V	2.41	2.41	0.00

### 3.3.1 Static Calculations

In addition with the relaxation calculations presented, static calculations were done to verify the optimal adsorption site for the metal adsorbate. Figures 5 to 8 show the results of these static calculations. If two sites have very close computed energy, the metal adatom may likely transfer to that adjacent site, until the minimum is reached. The

inset of Figs. 5 to 8 suggests a possible diffusion path of metal adatom towards the hollow site. In Fig. 5 (inset), the H site will most likely be preferred for the Co adatom from  $z > 3.0 \text{ \AA}$ . The atom may transfer to the V site at  $z = 2.5 \text{ \AA}$ , but it will again transfer to the hollow site at  $z = 1.5 \text{ \AA}$  and will stay there since the energy of other adsorption sites will be very far from the H site.

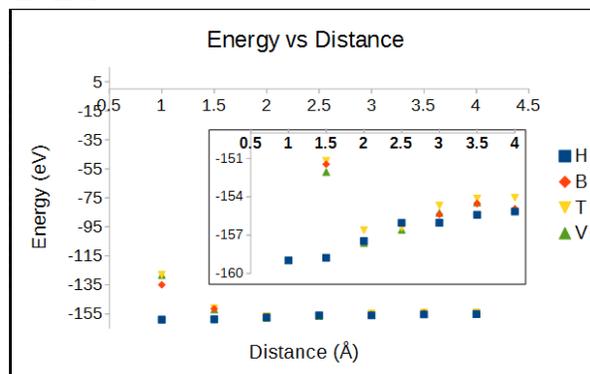


Fig. 5. Energy vs Distance above the adsorption site for Co.

The same can be seen for Ni (Fig. 6, inset). The Ni adatom may transfer sites at  $z > 1.00 \text{ \AA}$ , but will ultimately stay at H site since the energy of other sites are much higher than H at  $z = 1.00 \text{ \AA}$ .

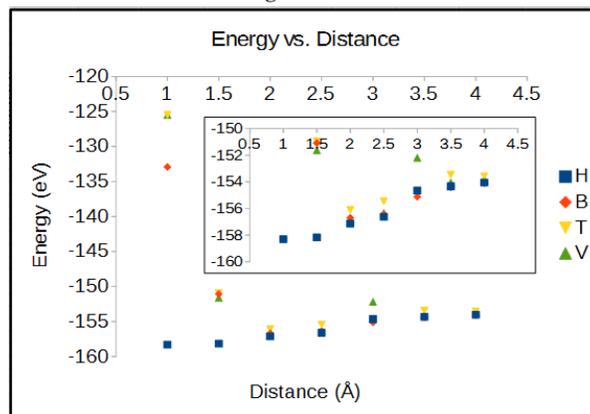


Fig. 6. Energy vs Distance above the adsorption site for Ni.

For the Pd and Pt adatoms (Fig.6-7, inset), it may move around the B, H, and V sites, at  $z > 2.5 \text{ \AA}$ , but it will eventually transfer at the lowest point, which is the H site at  $z = 1.5 \text{ \AA}$ .

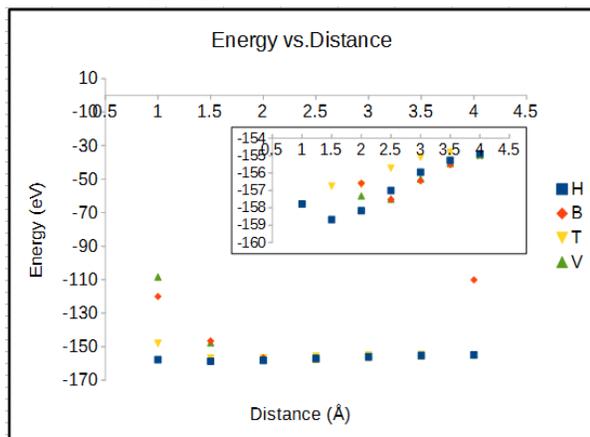


Fig. 7. Energy vs Distance above the adsorption site for Pd.

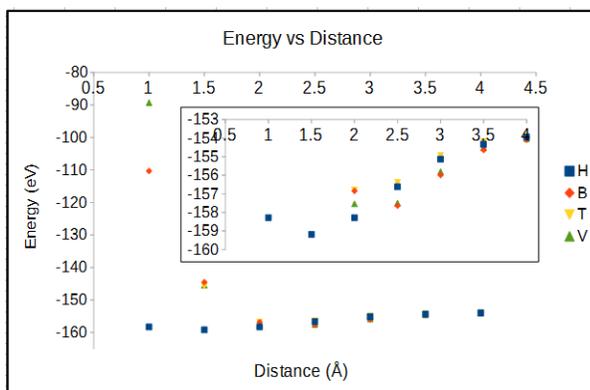


Fig. 8. Energy vs Distance above the adsorption site for Pt.

The static calculations confirmed the results of relaxation calculations, which show the H site as the most stable adsorption site for the four metal atoms. It also suggests a possible diffusion path in silicene as the adatom moves towards the H site. However, there is a possibility that metal adatom may settle at a minimum other than the ones shown in Figs. 5 to 8, since the  $z$  increments used are relatively large ( $1.0\text{\AA}$ ) and the true minima for the adatoms may have been missed. Also, for the static calculations of Co and Ni, some computed energies did not meet the convergence criteria of  $1 \times 10^{-4}$  eV. Some of the Co and Ni calculations obtained energy differences ( $dE$ ) of the orders  $10^{-1}$  eV to  $10^{-3}$  eV, which is higher than the typically accepted  $dE$  values. Hence, recalculation of these values are highly recommended.

### 3.4 H Interaction on Metal-Decorated Silicene

Calculations with the Pt-decorated silicene are currently ongoing. Initial ground state energy results encountered high  $dE$  values, which needs further recalculation. Afterwards, H interaction on silicene with other transition metal atoms will be done.

## 4. CONCLUSIONS

The adsorption site of Co, Ni, Pd, and Pt atoms on silicene was determined using DFT calculations. The relaxation calculations agreed with the previous study of Lin and Ni (2012) and showed the hollow site to be the most stable site for adsorption. Static calculations were also done and confirmed the relaxation results. The static calculations also suggested a possible diffusion path of metal adatom on silicene, but increasing the increment of  $z$  is recommended to ensure that the true energy minimum was not missed. Moreover, some of the static calculations for Co and Ni did not meet the energy convergence criterion and hence demand recalculation. Further studies on  $H_2$  interaction with metal-decorated silicene are currently ongoing.

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